A thresholding algorithm for hyperbolic mean curvature flow

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反応拡散現象にみられる境界層とその周辺の数理 2014年11月28日

Outline

1. Curvature dependent motions

 $\star \quad v = -\kappa \quad \text{(Mean curvature flow)}$

 $\star \quad \frac{dv}{dt} = -\kappa \quad \text{(Hyperbolic MCF)}$

(see J. Math. Pures Appl., P. LeFloch and K. Smoczyk 2008)

2. An approximation method for HMCF

 \star The original BMO (mean curvature flow)

 \star The hyperbolic BMO

+ Justification of our thresholding algorithm.

- 3. Application of the numerical method.
 - \star Hyperbolic mean curvature flow (HMCF)
 - \star Multiphase HMCF
 - \star Volume preserving HMCF, contact angles, numerical tests

Interfacial dynamics

What does curvature-driven acceleration look like?



Mean Curvature Flow

$$v = -\kappa$$



Hyperbolic Mean Curvature Flow

$$\frac{dv}{dt} = -\kappa$$

Physical interpretation (membrane motion)

Wave equation

 $u_{tt} = u_{xx}$

 κ

Nonlinear wave equation

$$0 \qquad 4 \qquad \qquad u_{tt} = \operatorname{div}\left(\frac{\nabla u}{\sqrt{1 + |\nabla u|^2}}\right)$$

Hyperbolic mean curvature flow

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Computational Tools



Computational Tools



The scalar BMO algorithm (1992) (Bence-Merriman-Osher)

Evolve by mean curvature flow $v = -\kappa$ $0 < \Delta t \ll 1$

$$u_t = \Delta u \qquad \text{in } (0, \Delta t) \times \Omega$$

$$\frac{\partial u}{\partial \nu} = 0 \qquad \text{on } (0, \Delta t) \times \partial \Omega$$

$$u(t = 0, x) = \chi_E \qquad \text{in } \Omega$$



Reformulated multiphase (vector-type) BMO algorithm:



Curvature Dependent Motions



mean curvature flow

Volume constrained motions



Vector-type minimizing movement: $\mathbf{u}, \mathbf{u}_{n-1} \in H^1(\Omega; \mathbf{R}^{N-1})$

$$\mathcal{F}_{n}(\mathbf{u}) = \int_{\Omega} \left(\frac{|\mathbf{u} - \mathbf{u}_{n-1}|^{2}}{2h} + \frac{|\nabla \mathbf{u}|^{2}}{2} \right) dx \quad \begin{array}{l} \text{Heat Type} \\ \tilde{\epsilon} > 0 \\ + \sum_{k=1}^{N-1} \frac{1}{\tilde{\epsilon}} |V_{k} - meas(E_{k}^{n})|^{2} & \begin{array}{l} \text{Penalties} \\ \end{array} \\ \mathbf{E}_{k}^{n} = \{x \in \mathbf{R}^{m} : \mathbf{u}(x) \cdot \mathbf{p}_{k} = \max_{i=1,\dots,N} \mathbf{u}(x) \cdot \mathbf{p}_{i}\} \end{array}$$

Multi-phase volume preserving



Multi-phase volume preserving



Hyperbolic Mean Curvature Flow

$$\begin{cases} \alpha''(t,s) = -\kappa(s) & (t,s) \in (0,T) \times [0,1) \\ \alpha'(t=0,s) = v_0(s) & s \in [0,1) \\ \alpha(t=0,s) = \gamma(s) & s \in [0,1) \end{cases}$$

Hyperbolic Mean Curvature Flow

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The geometrical setting



The idea (= follow the level set) (=



Contours of the level set function:

$$u(t, \mathbf{x}(t)) = c$$
 $(c \in \mathbf{R})$

Encode the interface in the zero level set: $u(t, \mathbf{x}(t)) = 0$

Follow the path of a particle along the level set $\frac{d}{dt}u(t, \boldsymbol{x}(t)) = 0 \implies u_t(t, x(t)) + \nabla u(t, x) \cdot \dot{\boldsymbol{x}}(t) = 0$ The idea (= follow the level set) (=



The velocity has a normal and tangental component: $\dot{\boldsymbol{x}}(t) = (V_n \boldsymbol{n} + V_T \boldsymbol{T}) \quad u_t + \nabla u \cdot (V_n \boldsymbol{n} + V_T \boldsymbol{T}) = 0$

Rewrite using the normal vector: $u_t + |\nabla u| \mathbf{n} \cdot (V_n \mathbf{n} + V_T \mathbf{T}) = 0$ $\mathbf{n} = \frac{\nabla u}{|\nabla u|}$

Using orthogonality relations: $u_t = -V_n |\nabla u|$ (the level set equation) The idea (= follow the level set) (=



The interfacial acceleration: $\frac{\partial}{\partial t} \left[\frac{u_t}{|\nabla u|} \right] = -\dot{V}_n \implies \frac{u_{tt} |\nabla u| - |\nabla u|_t u_t}{|\nabla u|^2} = -\dot{V}_n$ $\Gamma_t = \{x \in \Omega : u(t, x) = 0\}$ The signed distance function

$$d(t,x) = \begin{cases} \inf_{y \in \Gamma_t} ||x - y|| & \text{if } x \in \{u(t,x) > 0\} \\ -\inf_{y \in \Gamma_t} ||x - y|| & \text{otherwise} \end{cases}$$

satisfies the Eikonal equation:

$$|\nabla d| = 1$$
 a.e. x

Choose the evolution: $u_{tt} = \Delta u$

$$-\dot{V}_{n} = \frac{u_{tt} \left|\nabla u\right| - \left|\nabla u\right|_{t} u_{t}}{\left|\nabla u\right|^{2}} = \frac{\Delta u \left|\nabla u\right| - \left|\nabla u\right|_{t} u_{t}}{\left|\nabla u\right|^{2}} = \frac{\left(\nabla \cdot \nabla u\right) \left|\nabla u\right| - \left|\nabla u\right|_{t} u_{t}}{\left|\nabla u\right|^{2}}$$

$$=\frac{\left(\nabla\cdot\left(\left|\nabla u\right|\boldsymbol{n}\right)\right)\left|\nabla u\right|-\left|\nabla u\right|_{t}u_{t}}{\left|\nabla u\right|^{2}}\implies-\dot{V}_{\boldsymbol{n}}=\nabla\cdot\boldsymbol{n}=\kappa$$

The hyperbolic BMO algorithm



Theorem: Let $2d_k - d_{k-1}$ evolve by the wave equation (*), for a time t > 0. Then the zero level set of the solution evolves with normal acceleration: $v' = -\kappa + O(t)$ (and velocity $v(t) = v(0) - t\kappa + O(t^2)$).

$$\begin{cases} u_{tt} = \Delta u & \text{in } (0, \Delta t) \times \Omega \\ \frac{\partial u}{\partial \nu} = 0 & \text{on } (0, \Delta t) \times \partial \Omega \\ u(t = 0, x) = 2d_k - d_{k-1} & \text{in } \Omega \\ u_t(t = 0, x) = 0 & \text{in } \Omega \end{cases}$$
(*)

Proof: (sketch, 2D) $x = (x_1, x_2)$

The solution to the wave equation with initial condition u_0 and initial velocity v_0 can be written:

$$u(x,t) = \frac{1}{2\pi t} \int_{B(x,t)} \frac{u_0(y) + tv_0(y) + \nabla u_0(y) \cdot (y-x)}{\sqrt{t^2 - |y-x|^2}} \, dy.$$

Consider an initial condition that is the sum of two signed distance functions:

$$u_0 = ad_0 + bd_{-1}, v_0 = 0 \ (a, b, \in \mathbf{R})$$

A Hyperbolic BMO Algorithm: the justification

Theorem: (LeFloch et al., J. Math. Pures Appl., 2008, 90) If the velocity field is normal to the hypersurface at time zero, then it is normal throughout the evolution.

Lemma: (Essedoglu et al., J. Comp. Phys., 2010, 229) Let f(x) be a smooth function whose graph (x, f(x)) describes the interface in a neighborhood of the origin. Then signed distance function d(x, f(x)) has the following Taylor expansion at x = 0:



Distance functions

Def. Given $A \subset \mathbf{R}^N$, the *distance function* from a point x to A is defined as:

$$d_A(x) = \inf_{y \in A} |y - x|.$$

Remark. Suppose $\partial \Omega$ is the boundary of a set Ω . Then $x \in \partial \Omega \implies d_{\partial \Omega}(x) = 0$

Thm. The map $x \to d_A(x)$ is uniformly Lipschitz continuous in \mathbf{R}^N :

$$\forall x, y \in \mathbf{R}^N, \quad |d_A(y) - d_A(x)| \le |y - x|.$$

proof. For all $z \in A$ and $x, y \in \mathbf{R}^N$

$$|z - y| \le |z - x| + |y - x|$$

$$d_A(y) = \inf_{z \in A} |z - y| \le \inf_{z \in A} |z - x| + |y - x| = d_A(x) + |y - x|.$$

Distance functions

Thm. The distance function satisfies the Eikonal equation:

$$\nabla d_{\partial\Omega}(x) = 1$$
 a.e. in \mathbf{R}^N .

proof.

- 1. For a given location x, let x_c denote the point on $\partial \Omega$ that is closest to x.
- 2. Consider the line segment connecting x to x_c .
- 3. Note that x_c is the closest point on $\partial \Omega$ to every point along this segment.
- 4. $-\nabla d_{\partial\Omega}(x)$ gives the direction of steepest descent.

Distance functions

Since distance functions have "kinks" on their zero level set, signed distance functions are often used when we need to compute derivatives on $\partial\Omega$.

$$signd_{\partial\Omega}(x) = \begin{cases} d(x,\partial\Omega) & x \in \Omega, \\ -d(x,\partial\Omega) & x \in \mathbf{R}^N \setminus \Omega. \end{cases}$$
$$x^* \in \partial\Omega$$
$$x^* \in \partial\Omega$$
$$d_{\partial\Omega}(x) & signd_{\partial\Omega}(x) \end{cases}$$

Thm. The signed distance function d(x, f(x)) has the following Taylor expansion at x = 0:

$$d(x,y) = y + \frac{1}{2}\kappa(0)x^{2} + \frac{1}{6}\kappa_{x}(0)x^{3} - \frac{1}{2}\kappa^{2}(0)x^{2}y + \frac{1}{24}(\kappa_{xx}(0) - 3\kappa^{3}(0))x^{4} - \frac{1}{2}\kappa(0)\kappa_{x}(0)x^{3}y + \frac{1}{2}\kappa^{3}(0)x^{2}y^{2} + O(|\boldsymbol{x}|^{5})$$

proof. (Essedoglu et al.) This follows directly from the following four lemmas.



Lemma 1. For sufficiently small y, we have

d(0,y) = y.

$$d_y(0,y) = 1$$
 $\frac{\partial^k}{\partial y^k} d(0,y) = 0$ (for $k = 2, 3, 4, ...$) and $d_x(0,y) = 0$

proof.

From the previous discussion, we already have d(0, y) = y. The partial derivatives with respect to y yield the first two expressions, and the last expression then follows from the eikonal equation.

Lemma 2. The following hold:

$$\frac{\partial^k}{\partial y^k} d_x(0, y) = 0 \quad (k = 1, 2, \dots)$$

for all sufficiently small (x, y).

proof.

Let $A(x, y) = d_x^2(x, y) + d_y^2(x, y)$. The eikonal equation implies A(x, y) = 1. Differentiating with respect to x and y:

$$\frac{1}{2}\frac{\partial}{\partial x}A(x,y) = d_x(x,y)d_{xx}(x,y) + d_y(x,y)d_{xy}(x,y) = 0 \qquad (\bigstar)$$
$$\frac{1}{2}\frac{\partial}{\partial y}A(x,y) = d_x(x,y)d_{xy}(x,y) + d_y(x,y)d_{yy}(x,y) = 0.$$

Evaluating at x = 0 and using the result of lemma 1, one has

$$d_{xy}(0, y) = 0$$
 (for all small enough y .)

Differentiation with respect to y yields the claim.

Lemma 3. The following hold:

 $d_{xx}(0,0) = \kappa(0),$ $d_{xxy}(0,0) = -\kappa^2(0),$ $d_{xxx}(0,0) = \kappa_x(0).$

proof.

Along the interface, $d_{xx}(x, f(x)) + d_{yy}(x, f(x)) = \kappa(x)$. Evaluating at x = 0 and using lemma 1 yields the first claim. To obtain the second claim, we differentiate (\star) with respect to x again:

$$\frac{1}{2}A_{xx}(x,y) = d_{xx}^2 + d_x d_{xxx} + d_{xy}^2 + d_y d_{xxy} = 0.$$

Evaluating at (x, y) = (0, 0) gives $d_{xxy}(0, 0) = -\kappa^2(0)$ (where we have used the previous results).

To obtain the last claim, we differentiate the expression for the Laplacian of the distance function with respect to x:

$$\frac{\partial}{\partial x} \left(d_{xx}(x, f(x)) + d_{yy}(x, f(x)) \right) = \kappa_x(x)$$

$$\parallel$$

$$d_{xxx}(x, f(x)) + d_{xxy}(x, f(x))f'(x) + d_{yyx}(x, f(x)) + d_{yyy}(x, f(x))f'(x)$$

$$(f(0) = 0, f'(0) = 0, f''(0) = -\kappa(0))$$

$$\Longrightarrow$$

 $d_{xxx}(0,0) + d_{xxy}(0,0)f'(0) + d_{yyx}(0,0) + d_{yyy}(0,0)f'(0) = -\kappa_x(0)$

Lemma 4. The following hold:

$$d_{xxxy}(0,0) = -3\kappa(0)\kappa_x(0),$$

$$d_{xxyy}(0,0) = 2\kappa^3(0),$$

$$d_{xxxx}(0,0) = \kappa_x x(0) - 3\kappa(0)^3.$$

proof.

Differentiating with respect to x:

$$\frac{1}{2}A_{xxx}(x,y) = \frac{\partial}{\partial x} \left(d_{xx}^2 + d_x d_{xxx} + d_{xy}^2 + d_y d_{xxy} \right) = 0$$
$$= 3d_{xx}d_{xxx} + 3d_{xy}d_{xxy} + d_x d_{xxxx} + d_y d_{xxyx}$$
Evaluated at $(x,y) = (0,0)$ yields

$$3\kappa(0)\kappa_x(0) + 0 + 0 + d_{xxyx}(0,0) = 0.$$

Differentiating with respect to y:

$$\frac{1}{2}A_{xxy}(x,y) = \frac{\partial}{\partial y} \left(d_{xx}^2 + d_x d_{xxx} + d_{xy}^2 + d_y d_{xxy} \right) = 0$$

 $= 2d_{xx}d_{xxy} + d_{xy}d_{xxx} + d_xd_{xxxy} + 2d_{xy}d_{xyy} + d_{yy}d_{xxy} + d_yd_{xxyy}$

Evaluating at (x, y) = (0, 0):

 $0 = -2\kappa(0)\kappa^2(0) + 0 + 0 + 0 + 0 + d_{xxyy}(0,0).$

$$2\kappa^3(0) = d_{xxyy}(0,0).$$

Differentiating a previous result with respect to x again, we obtain:

$$\frac{\partial}{\partial x} \left(d_{xxx}(0,0) + d_{xxy}(0,0) f'(0) + d_{yyx}(0,0) + d_{yyy}(0,0) f'(0) \right) = -\kappa_{xx}(0)$$
$$\Longrightarrow$$

 $d_{xxxx}(x, f(x)) + d_{xxxy}(x, f(x))f'(x) + (d_{xxxy}(x, f(x)) + d_{xxyy}(x, f(x))f'(x))f'(x) + f''(x)d_{xxy}(x, f(x)) + d_{yyyx}(x, f(x)) + d_{yyyx}(x, f(x))f'(x) + (d_{yyyx}(x, f(x) + d_{yyyy}(x, f(x)))f'(x))f'(x) + f''(x)d_{yyy}(x, f(x)) = \kappa_{xx}(x)$

Evaluated at x = 0, one obtains:

$$d_{xxxx}(0,0) + 0 + 0 + \kappa^3(0) + 2\kappa^3(0) + 0 + 0 + 0 = \kappa_{xx}(0).$$

 $d_{xxxx}(0,0) = \kappa_{xx}(0) - 3\kappa^3(0)$

The hyperbolic BMO algorithm







Multiphase motions



Properties of wave propagation gives a multiphase algorithm

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The reformulated multiphase hyperbolic BMO algorithm



Numerical behavior of our approximation method

multiphase hyperbolic mean curvature flow



Numerical behavior of our approximation method

multiphase hyperbolic mean curvature flow



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Uniform Energy Estimates on the Minimizing Movement

Vector-type minimizing movement: $u, u_{n-1}, u_{n-2} \in H^1(\Omega; \mathbb{R}^{N-1})$

$$\mathcal{F}_{n}(\boldsymbol{u}) = \int_{\Omega} \left(\frac{|\boldsymbol{u} - 2\boldsymbol{u}_{n-1} + \boldsymbol{u}_{n-2}|^{2}}{2h^{2}} + \frac{|\nabla \boldsymbol{u}|^{2}}{2} \right) d\boldsymbol{x} \quad \begin{array}{l} \text{Wave Type} \\ \tilde{\epsilon} > 0 \\ + \sum_{k=1}^{N-1} \frac{1}{\tilde{\epsilon}} \left| V_{k} - meas(E_{k}^{n}) \right|^{2} & \begin{array}{l} \text{Penalties} \end{array} \right.$$

If there is no penalty, the minimizing movement converges:

$$u_{t}^{h} \stackrel{*}{\rightharpoonup} u_{t}, \qquad \nabla \bar{u}^{h} \stackrel{*}{\rightharpoonup} \nabla u, \qquad \nabla u^{h} \stackrel{*}{\rightharpoonup} \nabla u \qquad (\text{weakly star in } L^{\infty}(0,T;L^{2}(\Omega)))$$

$$\bar{u}^{h} \rightarrow u, \qquad u^{h} \rightarrow u \qquad (\text{strongly in } L^{2}(Q_{T})).$$

$$\frac{u_{t}^{h}(t) - u_{t}^{h}(t-h)}{h} \stackrel{*}{\rightharpoonup} u_{tt}, \qquad (\text{weakly star in } L^{\infty}(0,T;L^{2}(\Omega)))$$

$$\bar{u}^{h}(t,x) \stackrel{u_{n}(x)}{\longrightarrow} \stackrel{u_{n+1}(x)}{t} \qquad u^{h}(t,x)$$

$$\tilde{u}^{h}(t,x) \stackrel{u_{n}(x)}{\longrightarrow} \stackrel{u_{n+1}(x)}{t} \qquad u^{h}(t,x)$$

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$$v = -\kappa + k$$

$$v' = -\kappa + \bar{v}$$



$$v' = -\kappa + \bar{v}$$



Properties of the approximation scheme

Multiphase Volume Preserving HMCF



Contact Angles



Error Analysis



Summary

We introduced a method for approximating motion by hyperbolic mean curvature flow (HMCF)

The method is a threshold-dynamical algorithm, of the BMO type

The level set formulation suggested that thresholding evolution by the wave equation should yield the desired dynamics

Using the explicit representation formulas of the wave equation, we showed that the thresholding process yields motion by hyperbolic mean curvature flow, with an error of order t.

Numerical investigations suggest that volume preservation should be possible as well

References

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2. P. G. LeFloch, K. Smoczyk. *The hyperbolic mean curvature flow*, Journal de Mathématiques Pures et Appliquées, Vol. 90, No. 6 (2008), pp. 591-614.

Thank you for your attention